

Numerical fracture experiments using nonlinear nonlocal models



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Joint work with
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Overview of the talk

- Nonlinear peridynamic model
- Well-posedness of nonlocal model
- Numerical results
- Ongoing and future work

Nonlinear peridynamic model

Let ϵ is the horizon, $B_\epsilon(\mathbf{x})$ ball of radius ϵ , and $\mathbf{u}(\mathbf{x})$ displacement of material point $\mathbf{x} \in D$. In this work, we consider linearized pairwise strain $S(\mathbf{y}, \mathbf{x}; \mathbf{u})$ given by

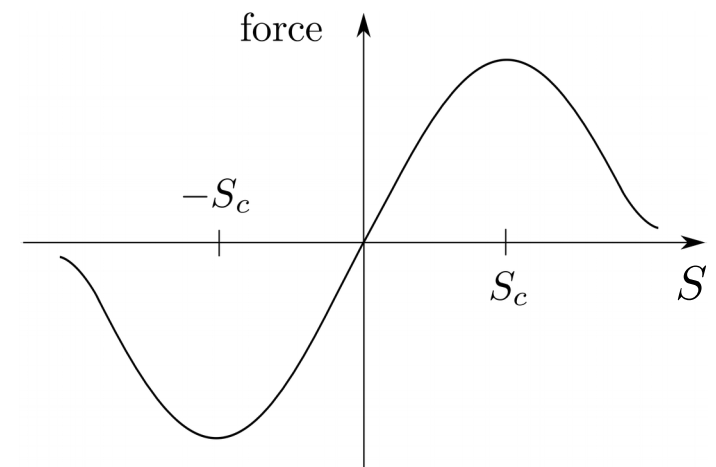
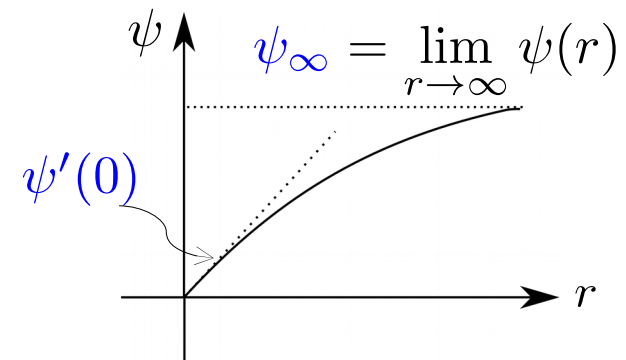
$$S(\mathbf{y}, \mathbf{x}; \mathbf{u}) = \frac{\mathbf{u}(\mathbf{y}) - \mathbf{u}(\mathbf{x})}{|\mathbf{y} - \mathbf{x}|} \cdot \frac{\mathbf{y} - \mathbf{x}}{|\mathbf{y} - \mathbf{x}|}.$$

Suppose $\hat{\mathbf{f}}^\epsilon(\mathbf{y}, \mathbf{x})$ denotes the force applied on \mathbf{x} from the neighboring point \mathbf{y} . Then total force at \mathbf{x} is given by

$$\mathbf{f}^\epsilon(\mathbf{x}) = \int_{B_\epsilon(\mathbf{x})} \hat{\mathbf{f}}^\epsilon(\mathbf{y}, \mathbf{x}) d\mathbf{y}$$

We consider pairwise force based on smooth and concave potential function $\psi^{1,2}$

$$\hat{\mathbf{f}}^\epsilon(\mathbf{y}, \mathbf{x}) = \frac{1}{\epsilon |B_\epsilon(\mathbf{0})|} \frac{\partial_S \psi(|\mathbf{y} - \mathbf{x}| S(\mathbf{y}, \mathbf{x})^2)}{|\mathbf{y} - \mathbf{x}|} \frac{\mathbf{y} - \mathbf{x}}{|\mathbf{y} - \mathbf{x}|}$$



[1] R. Lipton (2014) Dynamic brittle fracture as a small horizon limit of peridynamics. Journal of Elasticity, 117(1) 21-50.

[2] R. Lipton (2016) Cohesive dynamics and brittle fracture. Journal of Elasticity, 124(2), pp.143-191.

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Equation of motion

Equation of motion

$$\rho \ddot{\mathbf{u}}(\mathbf{x}, t) = \mathbf{f}^\epsilon(\mathbf{x}; \mathbf{u}(t)) + \mathbf{b}(\mathbf{x}, t), \quad \forall \mathbf{x} \in D, t \in [0, T]$$

Boundary condition

$$\mathbf{u}(\mathbf{x}, t) = \mathbf{g}(\mathbf{x}, t) \quad \forall \mathbf{x} \in D_u, t \in [0, T]$$

$$\mathbf{b}(\mathbf{x}, t) = \mathbf{f}_{ext}(\mathbf{x}, t) \quad \forall \mathbf{x} \in D_f, t \in [0, T]$$

$D_u, D_f \subset D$ are layer with finite volume (area in 2-d) on which displacement and external force, respectively, are specified. External force is applied in the form of body force.

Initial condition: $\mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0(\mathbf{x}), \dot{\mathbf{u}}(\mathbf{x}, 0) = \mathbf{v}_0(\mathbf{x})$ for all $\mathbf{x} \in D$.

Well-posedness of nonlinear peridynamic model 3

- Using the fact that nonlinear peridynamic force is bounded and Lipschitz continuous with respect to displacement field $\mathbf{u} \in L_0^2(D)$, the existence of solutions over any finite time domain $[0, T]$ is shown [1].
- To prove existence of solutions in more regular spaces, we introduce boundary function ω into Peridynamic force. $\omega(\mathbf{x}) = 1$ in the interior and smoothly decays to 0 as \mathbf{x} approaches boundary ∂D .
- To perform a priori error analysis of finite difference approximation, we consider Hölder space $C_0^{0,\gamma}(D)$, $\gamma \in (0, 1]$. In [2] we show existence of solutions in Hölder space $C_0^{0,\gamma}(D)$. In [3] we extend the results to state-based peridynamic models.
- For a priori error analysis of finite element approximation using continuous piecewise linear elements, we consider natural space $H^2(D) \cap H_0^1(D)$. In [4] we show existence of solutions in $H^2(D) \cap H_0^1(D)$.

[1] R. Lipton (2016) Cohesive dynamics and brittle fracture. *Journal of Elasticity*, 124(2), pp.143-191.

[2] P.K. Jha and R. Lipton (2018) Numerical analysis of nonlocal fracture models in Holder space. *SIAM Journal on Numerical Analysis*, 56(2), pp.906-941.

[3] P.K. Jha and R. Lipton (2019) Numerical convergence of finite difference approximations for state based peridynamic fracture models. *Computer Methods in Applied Mechanics and Engineering*, 351(1), 184 – 225.

[4] P.K. Jha and R. Lipton (2018) Finite element approximation of nonlocal fracture models. arXiv preprint arXiv:1710.07661. **Under review** in *Discrete and Continuous Dynamical Systems Series B*.

Finite difference approximation

We approximate peridynamic equation using **piecewise constant interpolation** and **central in time discretization**. Let \mathbf{u}_i^k denote the discrete displacement at mesh node \mathbf{x}_i and time $t^k = k\Delta t$. We consider following piecewise constant function

$$\mathbf{u}_h^k(\mathbf{x}) = \sum_{i, \mathbf{x}_i \in D} \mathbf{u}_i^k \chi_{U_i}(\mathbf{x})$$

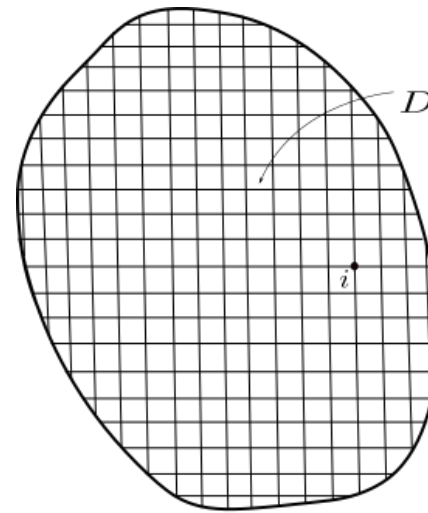
Discrete problem is

$$\frac{\mathbf{u}_h^{k+1} - 2\mathbf{u}_h^k + \mathbf{u}_h^{k-1}}{\Delta t^2} = \mathbf{f}_h^\epsilon(t^k) + \mathbf{b}_h^k,$$

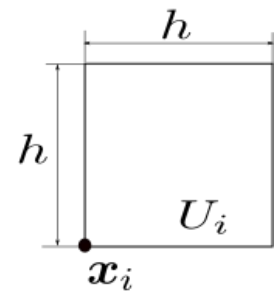
where

$$\mathbf{f}_h^\epsilon(\mathbf{x}, t^k) = \sum_{i, \mathbf{x}_i \in D} \mathbf{f}^\epsilon(\mathbf{x}_i, t^k) \chi_{U_i}(\mathbf{x}),$$

$$\mathbf{b}_h(\mathbf{x}, t^k) = \sum_{i, \mathbf{x}_i \in D} \mathbf{b}(\mathbf{x}_i, t^k) \chi_{U_i}(\mathbf{x})$$



(a)



(b)



Convergence of finite difference approximation

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Error at time step k is defined as: $E^k = \|\mathbf{u}_h^k - \mathbf{u}(t^k)\|$.

Theorem 1. *Let $\epsilon > 0$ be fixed. Let (\mathbf{u}, \mathbf{v}) be the solution of peridynamic equation. We assume $\mathbf{u}, \mathbf{v} \in C^2([0, T]; C^{0,\gamma}(D; \mathbb{R}^d))$. Then the finite difference scheme is consistent in both time and spatial discretization and converges to the exact solution uniformly in time with respect to the L^2 norm. If we assume the error at the initial step is zero then the error E^k at time t^k is bounded and satisfies*

$$\sup_{0 \leq k \leq T/\Delta t} E^k \leq O\left(C_t \Delta t + C_s \frac{h^\gamma}{\epsilon^2}\right),$$

where constant C_s and C_t are independent of h and Δt . Constants C_t, C_s depend on the ϵ and Hölder norm of the exact solution.

- [1] P.K. Jha and R. Lipton (2018) Numerical analysis of nonlocal fracture models in Holder space. SIAM Journal on Numerical Analysis, 56(2), pp.906-941.
 [2] P.K. Jha and R. Lipton (2019) Numerical convergence of finite difference approximations for state based peridynamic fracture models. Computer Methods in Applied Mechanics and Engineering, 351(1), 184 – 225.

Setting up peridynamic model

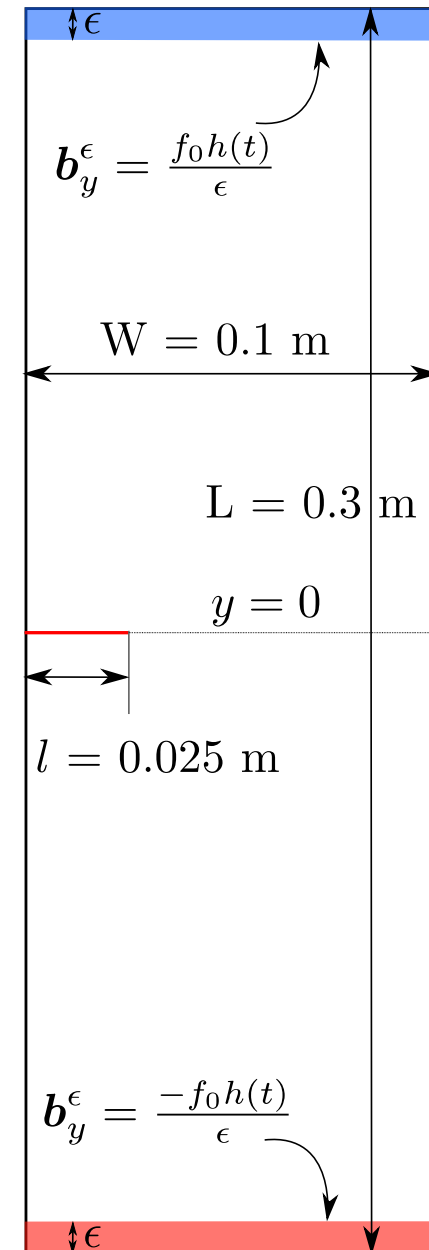
- Pairwise potential: $\psi(r) = c(1 - \exp[-\beta r^2])$
- Influence function: $J(r) = 1 - r$ for $0 \leq r < 1$ and $J(r) = 0$ for $r \geq 1$
- Critical strain: $S_c(\mathbf{y}, \mathbf{x}) = \frac{\pm \bar{r}}{\sqrt{|\mathbf{y} - \mathbf{x}|}}$, where \bar{r} is the inflection point of function ψ
- We fix $\rho = 1200 \text{ kg/m}^3$, bulk modulus $K = 3.24 \text{ GPa}$, critical energy release rate $G_c = 500 \text{ J/m}^{-2}$
- Using relation between nonlinear peridynamic model and linear elastic fracture mechanics¹, we find

$$c = 4712.4, \quad \beta = 1.7533 \times 10^8, \quad \bar{r} = \frac{1}{\sqrt{2\beta}} = 5.3402 \times 10^{-5}$$

Mode I crack propagation: Setup

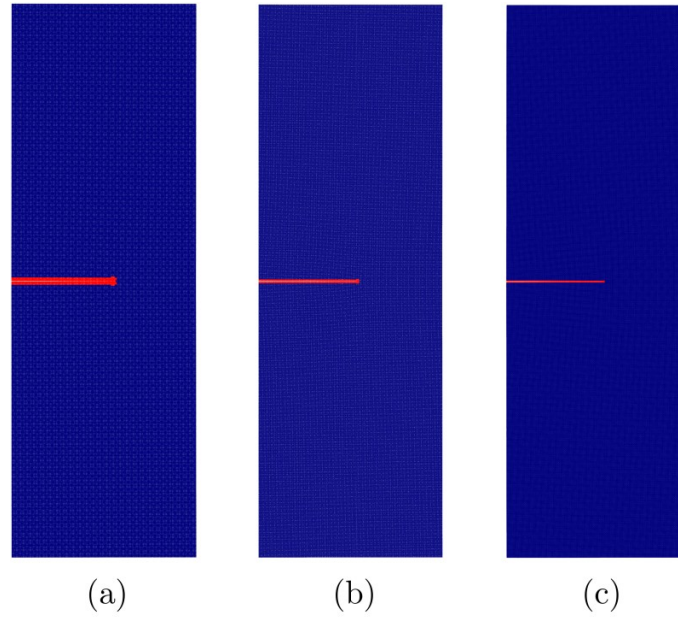
Goal: Localization of crack and convergence to classical fracture mechanics for simple mode-I crack propagation¹

- Final time $T = 560 \mu s$, time step $\Delta t = 0.02 \mu s$
- Uniform grid on square domain $D = [0, 0.1 \text{ m}] \times [-0.15 \text{ m}, 0.15 \text{ m}]$
- Experiment with three different horizons $\epsilon = 2.5, 1.25, 0.625 \text{ mm}$
- Body force $\mathbf{b}^\epsilon(\mathbf{x}, t) = (0, f_0 h(t)/\epsilon)$ on top layer and $\mathbf{b}^\epsilon(\mathbf{x}, t) = (0, -f_0 h(t)/\epsilon)$ on bottom layer
- $h(t)$ is a step function such that $h(t) = t$ for $t \leq 350 \mu s$ and $h(t) = 1$ for $t > 350 \mu s$
- Mesh size is fixed by relation $h = \epsilon/4$

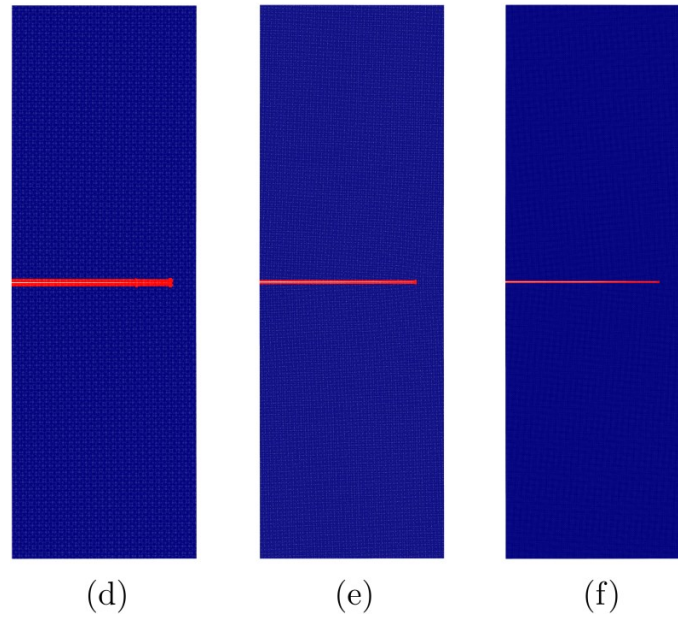


Localization of fracture zone

$t = 460 \mu s$

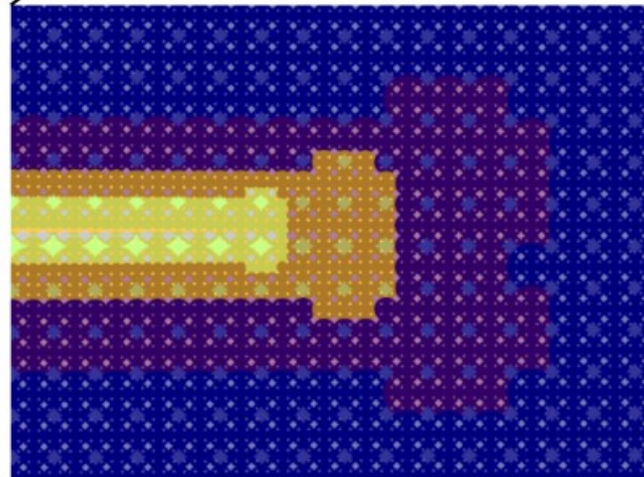
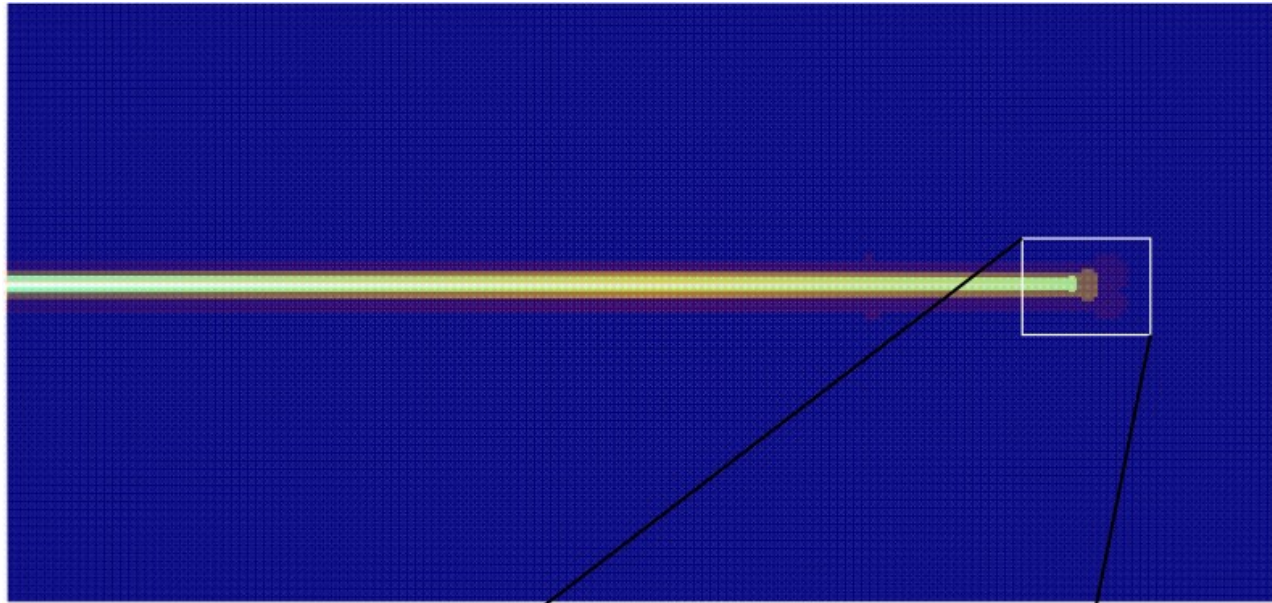


$t = 520 \mu s$

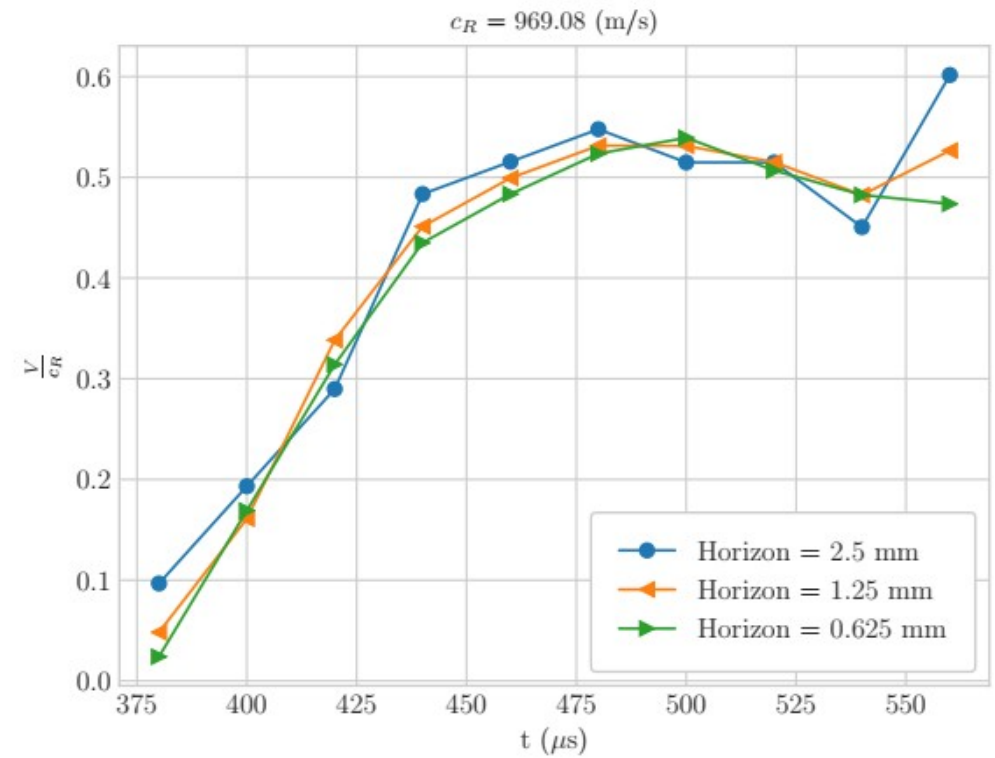
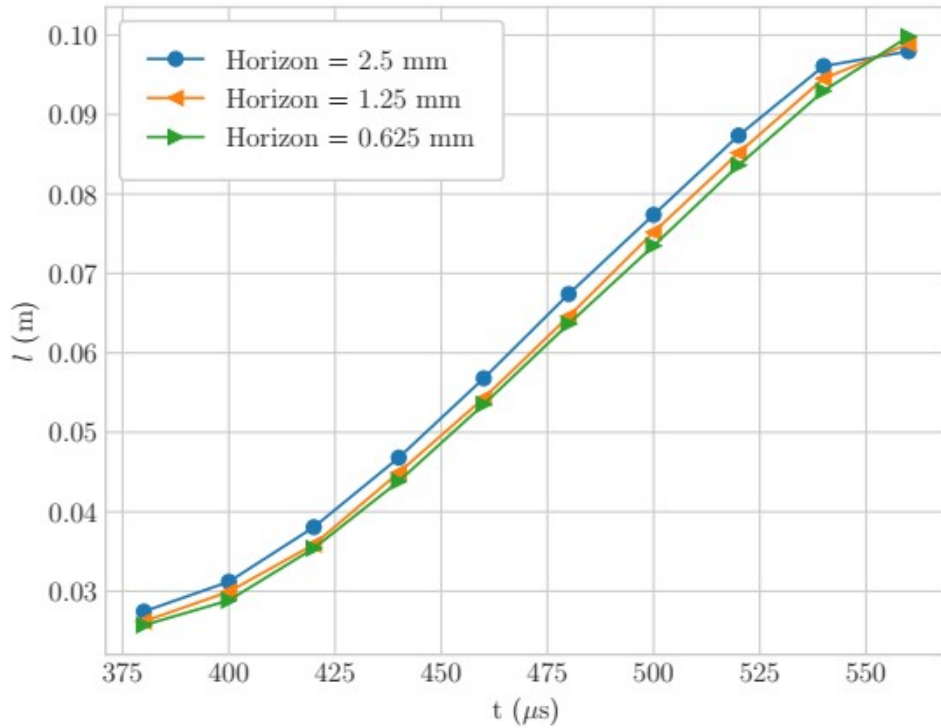


Localization of fracture zone

$t = 520 \mu s$



Crack tip location and velocity

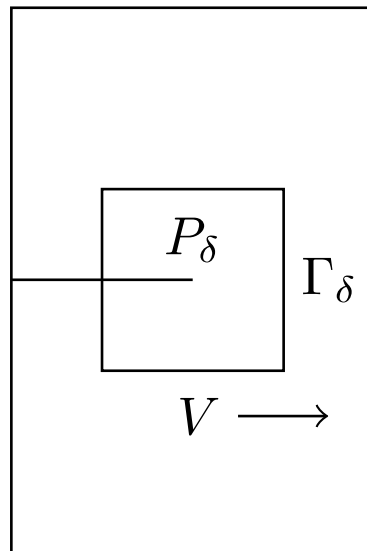


Energy into crack

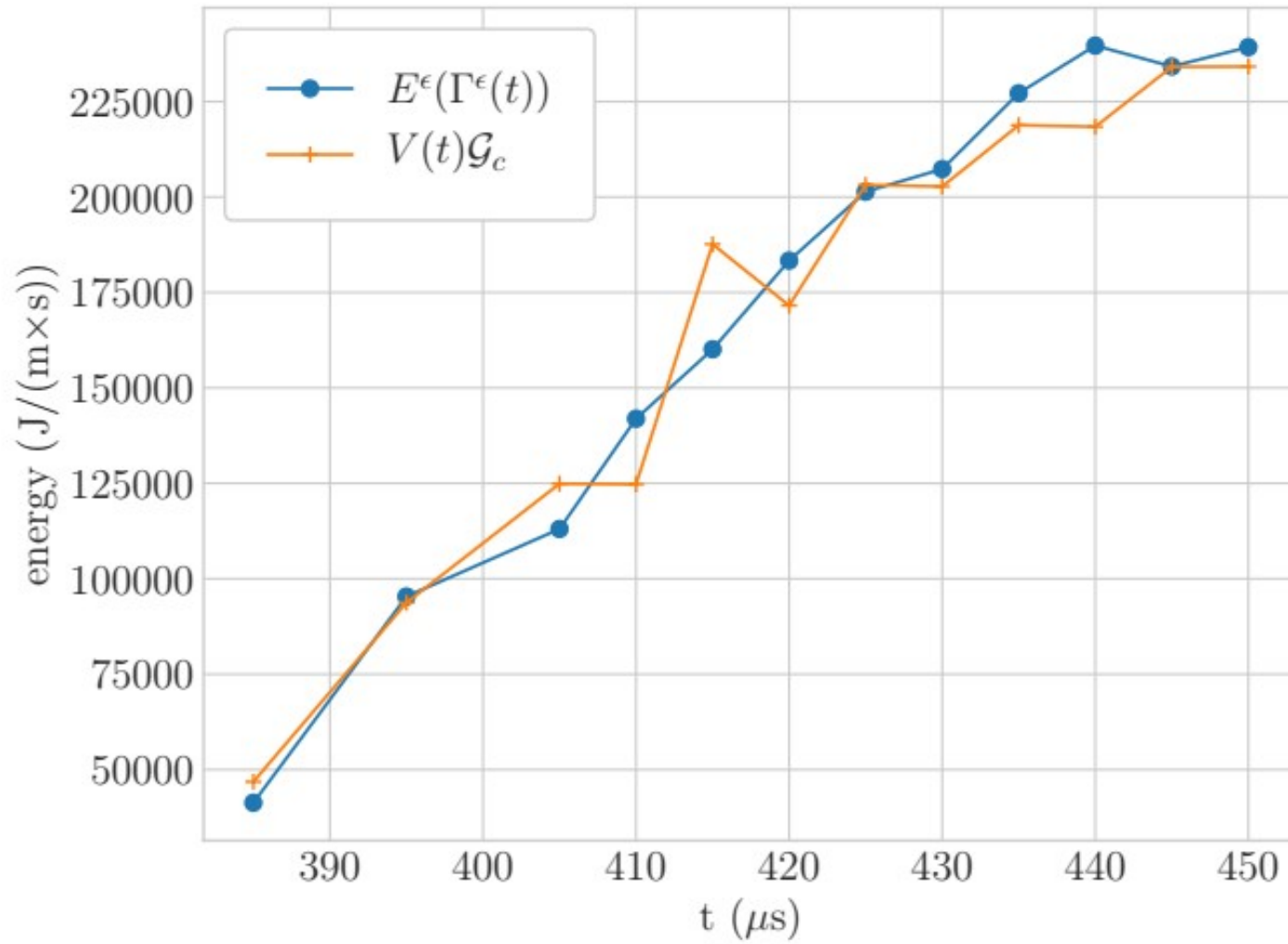
The energy associated to crack is given by¹

$$E(\Gamma_\delta(t)) = \frac{1}{|B_\epsilon(\mathbf{0})|} \int_{P_\delta^c(t)} \int_{P_\delta(t) \cap B_\epsilon(\mathbf{x})} \partial_S W(S(\mathbf{y}, \mathbf{x}; \mathbf{u}(t))) \frac{\mathbf{y} - \mathbf{x}}{|\mathbf{y} - \mathbf{x}|} \cdot (\dot{\mathbf{u}}(\mathbf{x}, t) + \dot{\mathbf{u}}(\mathbf{y}, t)) d\mathbf{y} d\mathbf{x}.$$

Here W is the peridynamic pairwise energy density. $P_\delta(t)$ is the rectangle domain with crack tip at its center. It is moving with tip. $P_\delta^c(t)$ is the complement of $P_\delta(t)$.



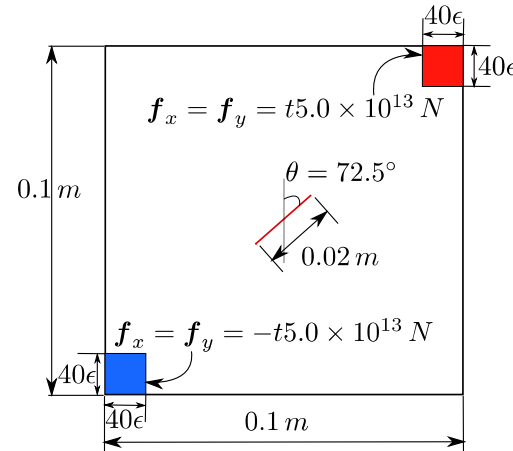
Energy into crack



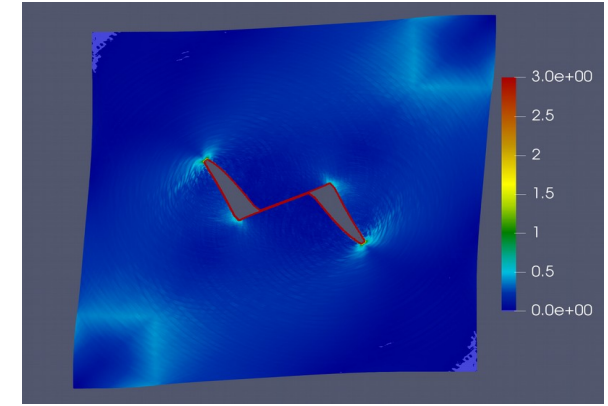
Mix mode crack propagation

Material properties are same as in the Mode-I problem. We set

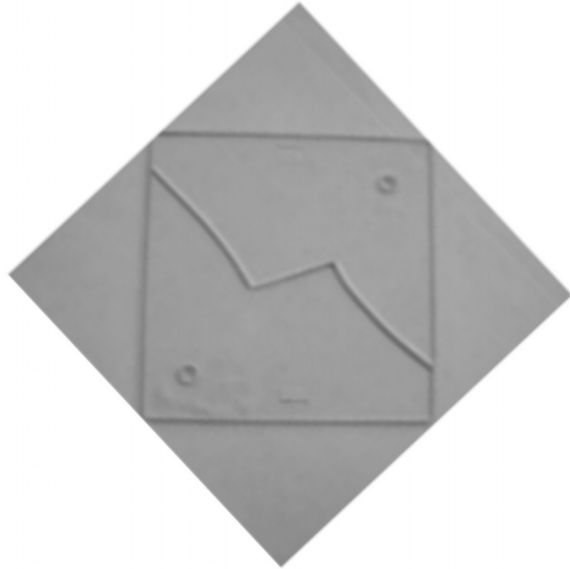
- Horizon $\epsilon = 0.5$ mm
- Mesh size $h = 0.125$ mm
- Final time $T = 140$ μ s
- Time step size $\Delta t = 0.004$ μ s



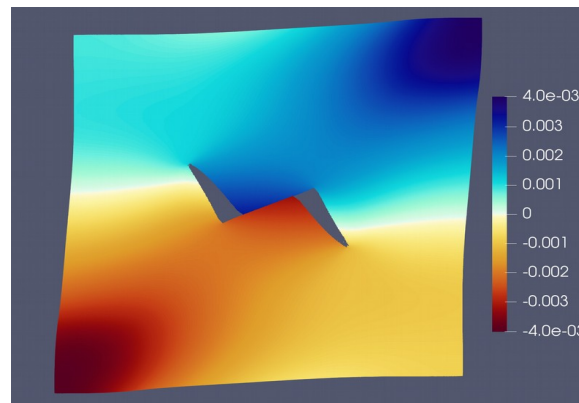
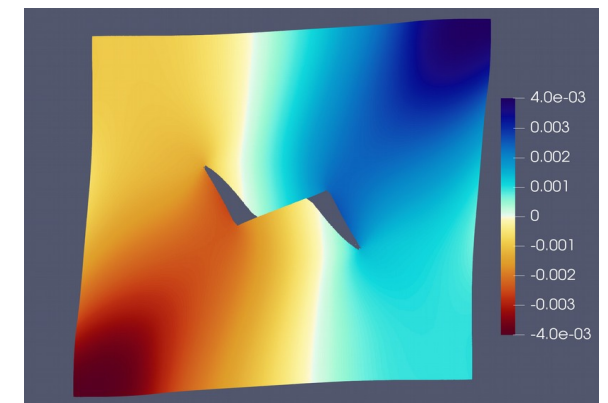
(a) Setup



(b) Damage profile



(e) Experiment result [2]


 (c) u_x plot

 (d) u_y plot

[1] R. Lipton, R. Lehoucq, & P.K. Jha (2019) Complex fracture nucleation and evolution with nonlocal elastodynamics. Journal of Peridynamics and Nonlocal Modeling. April 2019.

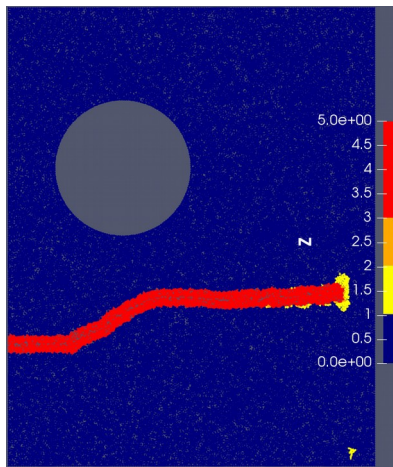
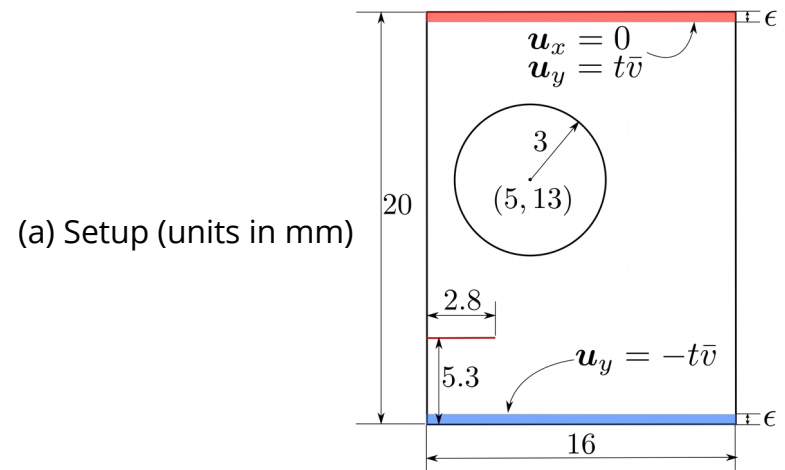
[2] M. R. Ayatollahi & M. R. M. Aliha (2009). Analysis of a new specimen for mixed mode fracture tests on brittle materials. Engineering Fracture Mechanics, 76(11), 1563-1573.

[3] E. Madenci et al (2018). A state-based peridynamic analysis in a finite element framework. Engineering Fracture Mechanics, 195, pp.104-128.

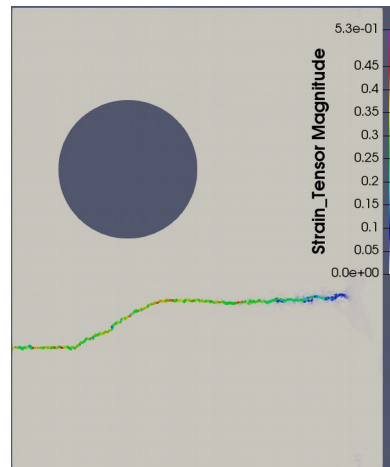
Crack-void interaction

Material properties are same as in the Mode-I problem. We set

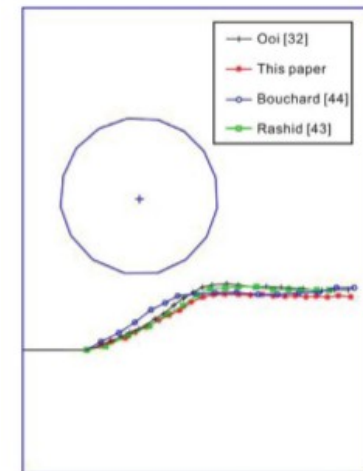
- Horizon $\epsilon = 0.4$ mm
- Mesh size $h = 0.1$ mm
- Final time $T = 800 \mu s$
- Time step size $\Delta t = 0.004 \mu s$



(b) Damage profile



(c) Magnitude of symmetric gradient of displacement



(d) Numerical experiment results using FEM, Boundary element method [2]

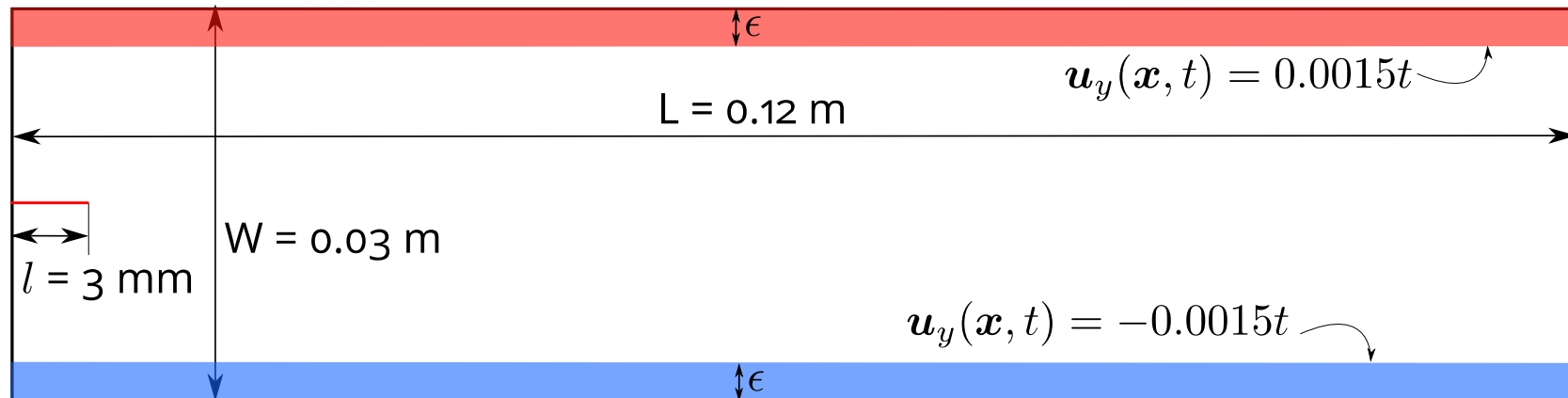
[1] P.K. Jha, P. Diehl & R. Lipton. Nodal finite element approximation of nonlocal fracture models. *In preparation*.

[2] S. Dai, C. Augarde, C. Du & D. Chen (2015). A fully automatic polygon scaled boundary finite element method for modelling crack propagation. *Engineering Fracture Mechanics*, 133, 163-178.

Wave reflection effect on crack velocity

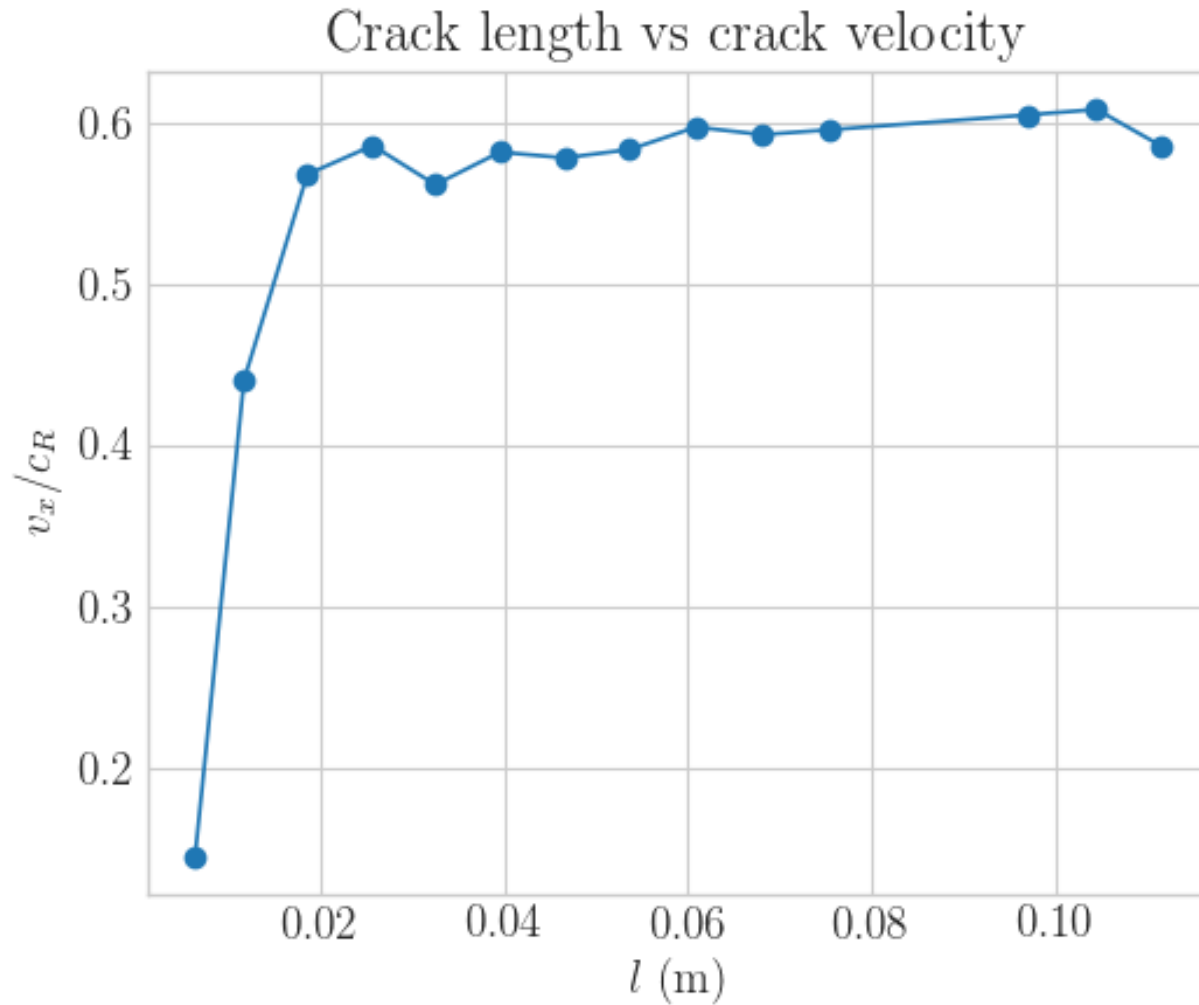
We consider a softer material with shear modulus $G = 35.2 \text{ kPa}$, density $\rho = 1011 \text{ kg/m}^3$, and critical energy release rate $G_c = 20 \text{ J/m}^{-2}$. Poisson ratio is fixed to $\mu = 0.25$. Domain is $D = [0, 0.12 \text{ m}] \times [0, 0.03 \text{ m}]$.

- Horizon $\epsilon = 0.6 \text{ mm}$, mesh size $h = 0.15 \text{ mm}$
- Time $T = 1.1 \text{ s}$, $\Delta t = 2.2 \mu\text{s}$



Wave reflection effect on crack velocity

- Max crack length = 0.12 m
- Rayleigh wave speed $c_R = 5.502$ m/s



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Ongoing and future works

- In [1] we show that the classical kinetic relation is embedded in peridynamics and we have $\lim_{\epsilon \rightarrow 0} J(t) = G_c$, where $J(t)$ is the nonlocal J-integral. In LEFM, the classical kinetic relation for the crack velocity is postulated. In contrast, we obtain the classical kinetic relation from the Peridynamics in the limit of vanishing nonlocality.
- Open source computational library for nonlocal modeling. This is a joint work with Patrick Diehl (LSU) and Robert Lipton (LSU).
- Study of granular material using nonlinear nonlocal model.



Thank you!